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Nous apprenons que la structure de $\text{LiCl} \cdot 2\text{en}$ vient d'être décrite brièvement par Brusset, Gillier-Pandraud & Delcroix (1966) dans une communication au Bulletin de la Société Chimique de France. Retenons que les distances interatomiques indiquées par les auteurs, bien que moins précises (écart-type moyen de 0,04 Å; $R=0,172$) sont comparables à celles que nous donnons ci-dessus, à l'exception toutefois des liaisons avec le lithium dont la position proposée est manifestement inexacte.*

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Two Independent Determinations of the Crystal and Molecular Structure of *trans*-Dichlorobis(dimethyl sulfoxide)palladium(II)

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The crystal and molecular structure of dichlorobis(dimethyl sulfoxide)palladium(II), $\text{PdCl}_2(\text{DMSO})_2$, has been determined from three-dimensional X-ray data in two independent studies. The complex crystallizes in the space group $P2_1/c$ of the monoclinic system with two molecules in a unit cell of dimensions (averages of the two studies) $a=6.460$, $b=9.380$, $c=9.555$ Å and $\beta=111.8^\circ$. Since the palladium atoms lie on crystallographic centers of symmetry the compound assumes a *trans* planar configuration. The dimethyl sulfoxide is found to be sulfur-bonded to the palladium atom in accordance with infrared spectral data. Anisotropic least-squares refinements led to residuals of 0.031 (with hydrogen atoms held fixed) and 0.118 (hydrogen atoms omitted) for the M.I.T. and Texas Christian studies, respectively. Some important dimensions (weighted averages) are: Pd-Cl 2.287 Å, Pd-S 2.298 Å, S-O 1.475 Å, S-C 1.778 Å, angle CSO 109.1° .

Introduction

The isolation of well-defined complexes containing dimethyl sulfoxide, $(\text{CH}_3)_2\text{SO}$ (hereafter abbreviated

DMSO), was first reported relatively recently (Lindquist & Einarsson, 1959; Cotton & Francis, 1960), though the powerful solvent properties of DMSO were generally recognized much earlier. A number of other

reports have appeared in which the formation of metal complexes of DMSO (Meek, Straub & Drago, 1960; Cotton & Francis, 1961) and other sulfoxides (Francis & Cotton, 1961) has been described. Attempts were made very early to establish the structures of these compounds indirectly on the basis of various physical measurements. One study of this sort, utilizing infrared spectra, led to the proposal (Cotton, Francis & Horrocks, 1960) that from shifts in the SO stretching frequency it could be concluded that in nearly all the known complexes the sulfoxide was attached to the metal (or other Lewis acid) through the oxygen atom. The exceptions postulated were complexes formed by Pd(II) and Pt(II), in which coordination by the sulfur atom was proposed.

The investigations described in this paper were undertaken to determine conclusively whether this postulate of sulfur bonding was correct and, in any event, to provide precise structure details upon which basis a close analysis of the metal to sulfoxide bonding could be founded. For various reasons the compound $\text{PdCl}_2(\text{DMSO})_2$ appeared to be very well suited to the purpose. A brief report of the M.I.T. study has been published earlier (Bennett, Cotton & Weaver, 1966).

Only after the two studies were complete, or very nearly so, was it realized that both had been carried out. It was therefore agreed, at the suggestion of Dr R. E. Marsh, that the two investigations should be published jointly, taking this somewhat unusual opportunity to compare the results of two entirely independent investigations which are comparable in the range and degree of refinement of the data. In the *Experimental* section, the two studies are described entirely separately [the M.I.T. work in part (1) and the Texas Christian work in part(2)]. Comparison and combining of results is done in the later sections.

Experimental (1)

Preparation

$\text{PdCl}_2(\text{DMSO})_2$ was prepared by the method of Cotton & Francis (1960) and crystals suitable for X-ray studies were obtained by slow diffusion of benzene into a solution of the complex in methylene chloride. The physical properties of the sample were in good agreement with those reported previously for the compound.

Collection and reduction of data

The crystals were examined by the precession method and found to belong to the uniquely determined space group $P2_1/c$ (C_{2h}^5 , no. 14) from the systematic absences $h0l$, $l=2n+1$, and $0k0$, $k=2n+1$. From precession photographs and measurements made with a General Electric Single Crystal Orienter the following unit-cell dimensions were found:

$$a = 6.450 \pm 0.005, b = 9.38 \pm 0.01, c = 9.54 \pm 0.01 \text{ \AA}, \beta = 111.8^\circ \pm 0.1^\circ (\lambda\text{Cu } K\alpha, 1.5418; \lambda\text{Mo } K\alpha_1, 0.70926 \text{ \AA}).$$

The measured density was found to be $2.07 \pm 0.05 \text{ g.cm}^{-3}$ by flotation in a mixture of methyl iodide and ethyl iodide; this compares favorably with a calculated density of 2.066 g.cm^{-3} for two $\text{PdCl}_2(\text{DMSO})_2$ molecules in the unit cell.

Intensity data were collected on a single crystal in the form of an orange needle of dimensions $0.1 \times 0.1 \times 0.4 \text{ mm}$ at room temperature. The crystal was mounted on the diffractometer with a^* in the nonstandard setting $B2_1/a$ parallel to the spindle or φ axis. A modified version of Furnas's moving-crystal, moving-counter method was used (Cotton & Elder, 1964); the mosaic spread of the crystal led to a choice of 2.66° in 2θ for the scan. The 836 independent reflections accessible in the spherical range $\theta_{\text{Mo}} \leq 24.4^\circ$, were collected at room temperature ($\sim 22^\circ\text{C}$) and corrected for the Lorentz and polarization factors. The profile of intensity versus φ was plotted for several $h00$ reflections, which occurred at $\chi = 90.0^\circ$, and found to show a maximum variation of $\pm 5\%$. Since at this orientation there was a maximum variation in the X-ray path and because the angular dependence of the transmission factor is small for the range over which the data were collected ($\mu r_{\text{max}} = 0.5$; $0 \leq \theta \leq 24.4^\circ$), absorption corrections were neglected. The data, indexed for $B2_1/a$, were transformed into the standard setting $P2_1/c$.

Solution and refinement of the structure

Since there are two molecules in the unit cell of space group $P2_1/c$, the Pd atoms are required to lie on centers of symmetry. From a three-dimensional Patterson function the sulfur and chlorine atoms were readily located, although not distinguished from each other. Two cycles of least-squares refinement phased on these three atoms led to a residual R_F [equation (1)] of 0.23%.

$$R_F = \frac{\sum |F_o - F_c|}{\sum |F_o|} \quad (1)$$

A difference Fourier synthesis was then calculated and the oxygen and carbon atoms were readily located. Further isotropic refinement phased on all these atoms resulted in a residual of 0.096%.

Three more cycles of least-squares refinement were attempted during which both positional and thermal parameters began to oscillate by more than an e.s.d. The 87 weak reflections for which counting statistics were unsatisfactory were omitted from further refinements. The atoms were reset at the centers of their oscillations and isotropic least-squares refinement using a unit weighting scheme converged to $R_F = 0.068$ in three cycles. A difference Fourier map calculated from the final isotropic structure factors showed evidence of anisotropic thermal motion. Four cycles of least-squares refinement with inclusion of anisotropic temperature factors of the form

$$\exp[-(\beta_{11}h^2 + \beta_{22}k^2 + \beta_{33}l^2 + 2\beta_{12}hk + 2\beta_{13}hl + 2\beta_{23}kl)]$$

for all atoms led to a residual of 0.035.

An analysis of the variation of $(F_{\text{obs}} - F_{\text{calc}})^2$ with F_{obs} revealed very little divergence from unit weights.

Refinement with an empirical weighting scheme produced no significant variation in either parameters ($<1.7\sigma$) or e.s.d.'s ($<10\%$) and unit weights were therefore retained.

A difference Fourier synthesis phased on the last cycle of refinement revealed six peaks around the two carbon atoms of 0.5 to 0.75 e.Å⁻³ at reasonable locations for hydrogen atoms. Since the standard deviation of the electron density of the difference map was estimated to lie between 0.07 e.Å⁻³ (Cruickshank, 1949) and 0.11 e.Å⁻³ (Lipson & Cochran, 1953), these peaks were considered to be real. Thus, six hydrogen atoms were included in four additional cycles of anisotropic refinement leading to a residual of 0.031 . The hydrogen atoms were assigned fractional coordinates corresponding to the positions obtained from the difference Fourier synthesis and anisotropic temperature parameters approximately 10% greater than those of the carbon to which they were attached. None of the hydrogen parameters were varied in the refinement.

In the final cycle of refinement the largest shift in the parameters was half of the standard deviation, an average shift being one-fourth of the standard deviation. A final difference Fourier map calculated from the parameters obtained in the last cycle of refinement exhibited no electron density maximum exceeding 0.5 e.Å⁻³. Atomic scattering factors used throughout for the neutral light atoms were those of Ibers (1962). The neutral Pd scattering factor was obtained from a tabulation by Cromer & Waber (1965) and was corrected for the real part of anomalous dispersion (Dauen & Templeton, 1955). The observed and calculated structure factors are listed in Table 1. The final coordinates and thermal parameters are given in Table 3.

The following programs for the IBM 7094 computer were used and are gratefully acknowledged.

1. W.G.Sly, D.P.Shoemaker and J.H.Van den Hende. Two- and three-dimensional crystallographic Fourier summation, ERFR-2 (1962).

2. D.P.Shoemaker. M.I.T. X-ray goniometer package, MIXG-2 (1962).

3. C.T.Prewitt. A full matrix crystallographic least squares program, SFLSQ3 (1962).

4. J.S.Wood. Molecular geometry with estimated standard deviations, MGEOM (1964).

5. D.P.Shoemaker. Crystallographic bond distance, bond angle and dihedral angle computer program, DISTAN (1963).

6. D.P.Shoemaker and R.C.Srivastava. Anisotropic temperature factor interpreting program, VIBELL (1963).

7. R.C.Elder. General data reduction and presentation programs, RAWRE, PDATA and PUBTAB (1964).

Experimental (2)

Preparation

Palladium(II) chloride obtained from Englehard Inds. Inc. was used without further purification. Di-

methyl sulfoxide obtained from Fisher Scientific Company was dried over calcium hydride and vacuum distilled before use.

trans-PdCl₂(DMSO)₂ was prepared by dissolving palladium(II) chloride in DMSO at 100°C and allowing the solution to cool slowly. Single crystals of the bright orange material, suitable for X-ray studies, were obtained upon recrystallization from a DMSO-acetone mixture, m.p. 203–204°C. The crystals showed a tendency to twin, but examination with a polarizing microscope permitted the selection of untwinned specimens.

Collection and reduction of data

A crystal $0.23 \times 0.13 \times 0.13$ mm was used for the collection of all X-ray data. Precession photographs indicated that the crystal belonged to the monoclinic system. Systematic extinctions $0k0$, $k \neq 2n$; $h0l$, $l \neq 2n$ placed the crystal in space group $P2_1/c$ (no. 14) with the b axis unique. The unit-cell dimensions were obtained from the $h0l$ Weissenberg photograph and a b -axis layer line photograph (λ Mo $K\alpha_1$, 0.70296 Å). The films were calibrated by superimposed sodium chloride powder lines. The unit-cell dimensions are $a = 6.48 \pm 0.01$, $b = 9.380 \pm 0.005$, $c = 9.562 \pm 0.004$ Å, $\beta = 111^\circ 49' \pm 10'$. The measured density, $d_0 = 2.037$ g.cm⁻³ (floatation method), indicated two molecules per unit cell; $d_c = 2.055$ g.cm⁻³.

Levels $h0l$ to $h10l$ were recorded by the equi-inclination Weissenberg technique. The intensities of the reflections were estimated visually with calibrated strips prepared from reflections of similar area. In this way, 1258 non-zero reflections were obtained. Absorption corrections were applied with the use of Bond's values with $\mu R = 0.187$ (*International Tables for X-ray Crystallography*, 1959).

Solution and refinement of structure

Since the space group is $P2_1/c$ with two molecules per unit cell, the palladium atoms are required to occupy special positions which are located at centers of symmetry. This requires like ligands to be *trans* to each other and a *trans* square-planar configuration was assumed. An electron density projection was calculated with the signs derived from the palladium atom positions. Positions corresponding to sulfur and chlorine atoms in a square-planar array were evident although they were indistinguishable at this stage. A series of successive Fourier syntheses was used to locate the lighter atoms and to assign the sulfur and chlorine positions definitely.

The atomic scattering factors used in the structure factor calculations were those computed from relativistic Dirac-Slater wave functions (Cromer & Waber, 1965). Anomalous dispersion corrections of the type $f' = f_0 + \Delta f' + i\Delta f''$ were applied to the scattering factors (*International Tables for X-ray Crystallography*, 1962). In all cases, the imaginary term was negligible.

Initial refinement was carried out through use of a difference electron density map. A full-matrix least-

squares technique was used for the final refinement. The quantity minimized was $\Sigma w(F_o - F_c)^2$. A few reflections that were obviously affected by extinction were omitted in the final calculations. Several cycles with isotropic temperature factors decreased R ($R = \Sigma(F_o - F_c)/\Sigma F_o$) to 0.137 and three cycles with the anisotropic factors gave the final value of $R = 0.118$.

Three different weighting schemes were used. The resulting parameters all agreed within the estimated standard deviations, and the results for the unit weighting scheme are presented for comparison with the M.I.T. calculations. The other weighting schemes used were $F_o < 4|F_{\min}|$, $w = 1/(16|F_{\min}|^2)$ and $w = 1.0$; $F_o > 4|F_{\min}|$, $w = 1/|F_{\text{obs}}|^2$ and $w = |F_{\min}|/4|F_{\text{obs}}|$. The only noticeable change for the latter two schemes was a trend toward equalizing the two S-C distances.

The final shifts in all parameters were well below their estimated standard deviations. Table 1 lists the observed and calculated structure factors while Tables 2 and 3 give the final positional and thermal parameters.

The following computer programs were used:

1. W.H. Watson and R.J. Williams. Data reduction program (IBM 1620).
2. R. Shiono. Structure factor calculations in X-ray crystallography (IBM 1620).
3. R. Shiono. General three-dimensional Fourier synthesis program (IBM 1620).

Table 2. Fractional coordinates of atoms

	x	y	z
Pd	0.0	0.0	0.0
S	0.0642 (3)	0.1981 (2)	-0.1200 (2)
Cl	0.0638 (7)	0.1969 (5)	-0.1202 (4)
O	0.3628 (3)	-0.0114 (2)	0.1647 (2)
C(1)	0.3616 (7)	-0.0125 (7)	0.1652 (5)
C(2)	0.0578 (10)	0.1655 (6)	-0.2729 (5)
	0.0626 (29)	0.1697 (19)	-0.2710 (15)
	0.3166 (13)	0.2885 (9)	-0.0167 (8)
	0.3122 (34)	0.2851 (28)	-0.0168 (23)
	-0.1347 (13)	0.3333 (8)	-0.1324 (8)
	-0.1461 (34)	0.3348 (22)	-0.1362 (20)

Table 3. Anisotropic thermal parameters ($\times 10^4$)

	β_{11}	β_{22}	β_{33}	β_{12}	β_{13}	β_{23}
Pd	129 (2)	51 (1)	52 (1)	4 (1)	17 (1)	2 (1)
	115 (4)	50 (6)	47 (2)	11 (3)	23 (2)	5 (2)
S	191 (5)	53 (2)	47 (2)	-3 (3)	36 (3)	2 (2)
	212 (11)	52 (8)	55 (4)	-1 (5)	48 (5)	-2 (3)
Cl	151 (5)	85 (2)	111 (2)	-2 (3)	11 (3)	16 (2)
	150 (10)	104 (10)	116 (6)	3 (7)	3 (6)	14 (5)
O	411 (21)	93 (7)	65 (6)	-20 (10)	85 (10)	-13 (6)
	471 (61)	121 (25)	76 (16)	-16 (30)	127 (26)	-24 (15)
C(1)	223 (22)	89 (10)	97 (10)	45 (13)	41 (13)	-1 (8)
	204 (53)	124 (35)	109 (25)	63 (34)	47 (29)	-5 (22)
C(2)	286 (23)	65 (9)	94 (10)	45 (12)	47 (13)	13 (8)
	278 (57)	62 (25)	90 (21)	42 (28)	89 (28)	4 (16)

4. M.E. Pippy. Calculation of bond lengths and bond angles (IBM 1620).

5. H.A. Levy and W.R. Busing. Full-matrix least-squares (IBM 7094, Texas A and M University).

Description of the structure

Molecular structure

Table 4 gives the molecular dimensions while Fig. 1 shows the molecular structure of $\text{PdCl}_2(\text{DMSO})_2$ projected along the a axis. Since the Pd atom lies on a crystallographic center of symmetry, the molecule is required to have a *trans* planar configuration.

Table 4. Intramolecular distances* (\AA) and angles ($^\circ$) together with their estimated standard deviations

	Determination (1)		Determination (2)	
	Distance	e.s.d.	Distance	e.s.d.
Pd-S	2.299	0.002	2.291	0.005
Pd-Cl	2.287	0.002	2.290	0.006
S-O	1.476	0.005	1.46	0.02
S-C(1)	1.780	0.008	1.75	0.03
S-C(2)	1.776	0.007	1.84	0.03
	Angle	e.s.d.	Angle	e.s.d.
S-Pd-Cl	93.00	0.06	93.3	0.2
Pd-S-C(1)	114.1	0.3	113.4	1.0
Pd-S-C(2)	110.5	0.3	110.0	0.9
Pd-S-O	112.4	0.2	115.0	0.9
O-S-C(1)	109.2	0.3	107.6	1.3
O-S-C(2)	109.1	0.3	107.7	1.2
C(1)-S-C(2)	100.9	0.3	102.3	1.3
C(1)-H(av.)	1.07	†		
C(2)-H(av.)	0.97	†		

* All distances uncorrected for thermal motion.

† H atom positions derived from electron difference map.

Refinement (1)

The Pd-Cl distance of 2.30 \AA is exactly that found previously in a norbornadienedichloropalladium(II) complex (Baenziger, Doyle & Carpenter, 1961) while the Pd-S distance of 2.30 \AA compares quite favorably with that of 2.33 \AA in $\text{Pd}(\text{SCN}_2\text{H}_4)_2\text{Cl}_2$ (Oi, Kawase, Nakatso & Kuroya, 1960). The two crystallographically independent S-C distances are equal within ex-

perimental error at $1.778 \pm 0.007 \text{ \AA}$, while the S–O distance is $1.475 \pm 0.005 \text{ \AA}$. The environment around the sulfur atom is nearly tetrahedral with the angles ranging from 100.9 to 114.0° . The average C–H distances are reasonably close to the value of 1.09 \AA accepted (Pauling, 1960) for an sp^3 hybridized carbon atom bonded to hydrogen.

Refinement (2)

The Pd–S and Pd–Cl bond lengths are experimentally identical, 2.290 ± 0.006 and $2.291 \pm 0.005 \text{ \AA}$ respectively, and the bond angles around the palladium atom are 87.3 and 92.7° . The sulfur atom in each DMSO group is in an approximate tetrahedral environment with angles ranging from 102.3 to 115.0° . The two independent C–S bond lengths are 1.84 ± 0.03 and $1.75 \pm 0.03 \text{ \AA}$ in the unit weighting scheme and 1.82 ± 0.03 and $1.76 \pm 0.03 \text{ \AA}$ for the others.

Intermolecular distances

All intermolecular contacts up to 4 \AA were calculated. The pertinent ones are listed in Table 5. None of these are excessively short and no further significance is attached to them.

Analysis of anisotropic motion

Root mean square amplitudes of the atomic vibrational ellipsoid axes for $\text{PdCl}_2 \cdot (\text{DMSO})_2$ are given in

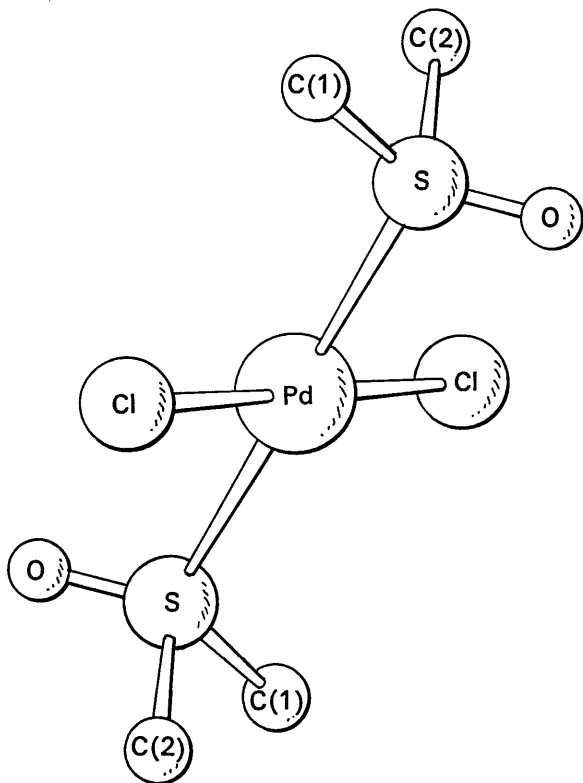


Fig. 1. The $\text{PdCl}_2(\text{DMSO})_2$ molecule viewed along the crystallographic a axis. The palladium atom lies on an inversion center.

Table 5. Intermolecular contacts in $\text{PdCl}_2(\text{DMSO})_2$

Atom of reference molecule	Atom in neighboring molecule	Distance
Cl	O ($x, \frac{1}{2}-y, \frac{1}{2}+z$)	3.95 \AA
	C(1) ($x, \frac{1}{2}-y, \frac{1}{2}+z$)	3.79
	C(1) ($\bar{x}, y-\frac{1}{2}, \frac{1}{2}-z$)	3.72
	C(2) ($x, \frac{1}{2}-y, \frac{1}{2}+z$)	3.50
S	O ($x, \frac{1}{2}-y, \frac{1}{2}+z$)	3.57
O	C(1) ($x, \frac{1}{2}-y, z-\frac{1}{2}$)	3.37
	C(2) ($x, \frac{1}{2}-y, z-\frac{1}{2}$)	3.20
	C(2) ($\bar{x}, y-\frac{1}{2}, -\frac{1}{2}-z$)	3.33

Table 6 together with their direction cosines. These are referred to principal axes in a cartesian system in which x is parallel to \mathbf{a} , y parallel to \mathbf{b} and z parallel to \mathbf{c}^* . The anisotropy of the individual atoms is that expected. For example, the Cl minor axis is along the Pd–Cl bond direction, the major axis being perpendicular to the plane of the complex and to the Pd–Cl bond. Similarly the oxygen and two carbon atoms have their minor axes in the S–O bond direction and S–C bond directions respectively. The palladium atom vibrates in such a way that the minor axis lies roughly along the Pd–Cl bond direction. The major axis lies slightly below the plane of the complex and bisects the S–Pd–Cl angle.

Comparison of results

A comparison of the two independent sets of parameters and their estimated standard deviations emphasizes the main factors in the accuracy of crystal structure determinations. The two main features affecting the precision of the determined values are (a) accuracy of data measurement, and (b) the number of observations. Clearly in this case determination (1), using diffractometer data, gains from criterion (a) while determination (2), from film data, is better under criterion (b). In this example it would appear that the improved accuracy in the measurement of the intensities more than offsets the larger number of observations to the extent that the values for the estimated standard deviations of the positions in determination (1) are, in general, one third the values of those for determination (2). The agreement between the sets of parameters is satisfying with no discrepancy exceeding 3σ , σ being that in determination (2).

Cruickshank (1960) gives a simple formula for coordinate e.s.d.'s for data that are isotropic in three dimensions and have no atoms on special positions:

$$\sigma(x) \approx \frac{R}{\bar{S}} \left(\frac{N}{8p} \right)^{\frac{1}{2}} \text{ (centrosymmetric case)}$$

where

R = residual $\Sigma |\Delta F| / \Sigma |FOBS|$

\bar{S} = r.m.s. reciprocal radius for the planes observed

p = difference between the number of observations and the numbers of parameters

N = the number of atoms, of the type being considered, that are needed to give a scattering power equivalent to the total asymmetric unit at \bar{S} .

Although both data sets are approximately isotropic, clearly the formula is invalid since the palladium atom is on a special position. However, the ratio of the e.s.d.'s of the two determinations should be closely approximated by the expression:

$$\frac{\sigma_1(x)}{\sigma_2(x)} \simeq \frac{R_1 \bar{S}_2}{R_2 \bar{S}_1} \left(\frac{N_1 p_2}{N_2 p_1} \right)^{\frac{1}{2}},$$

which on substitution of the relevant values, $R_1 = 0.031$, $R_2 = 0.118$, $\bar{S}_1 = 0.90$, $\bar{S}_2 = 1.05$,* $p_1 = 744$, $p_2 = 1198$, $N_1 \simeq N_2$, gives $\sigma_1(x)/\sigma_2(x) \simeq 1/2.5$, which is in reasonable agreement with the observed results.

A semiquantitative treatment of the β_{ij} 's is more difficult because the lack of isotropy in data set (2) becomes more important. This feature is clearly shown by the disparity in e.s.d.'s for β_{22} 's and β_{33} 's in determination (2), where a straight comparison is reasonably valid, *i.e.* conversion to u_{ij} 's involves approximately the same constant. However, once again the dominant feature in the comparable formula for u_{ij} 's is R , and the lower e.s.d.'s come from determination (1) despite the large amount of high angle data in determination (2).

Discussion

The structure of this compound is of interest chiefly in two respects. First, and most obvious, it shows that the earlier interpretation of infrared spectra, in which binding of the DMSO to palladium is through sulfur is, indeed, correct. The general applicability of the infrared criteria proposed (Cotton, Francis & Horrocks, 1960) is thus strongly suggested.

Second, the detailed structure of DMSO in $\text{PdCl}_2(\text{DMSO})_2$ invites comparison with the molecular structure parameters of DMSO itself. Two reports of the DMSO structure, determined by X-ray crystallography have recently appeared (Viswamitra & Kannan, 1966; Thomas, Shoemaker & Eriks, 1966). We shall make comparisons with the molecular parameters of Thomas, Shoemaker & Eriks because their structure is (a) based on data collected nearer to room temperature ($+5^\circ$) than that of Viswamitra & Kannan (-60°), (b)

their data are more extensive and apparently better scaled, and (c) their refinement appears to have been more complete. Moreover, the thermal ellipsoids obtained in the M.I.T. study of $\text{PdCl}_2(\text{DMSO})_2$ and those of Thomas, Shoemaker & Eriks for DMSO are comparable in magnitude and shape. Since librational corrections to the bond lengths have not been applied in the present work, comparisons will be made with the uncorrected bond lengths in the DMSO structure.

The weighted mean S-C distance found in the present studies is $1.778(5) \text{ \AA}$. This is in good agreement with the values of $1.771(8)$ and $1.805(11)$ found by Thomas, Shoemaker & Eriks and the values of $1.80(1)$ and $1.81(1)$ found by Viswamitra & Kannan in DMSO itself, as well as with the S-C distances in the $(\text{CH}_3)_3\text{SO}^+$ ion (Zimmermann, Barlow & McCullough, 1963; Coulter, Gantzel & McCullough, 1963).

The S-O bond length is of particular interest. In resonance terms, the electron distribution can be expressed as a hybrid of the three canonical forms (I), (II) and (III), with (II) presumably predominating.

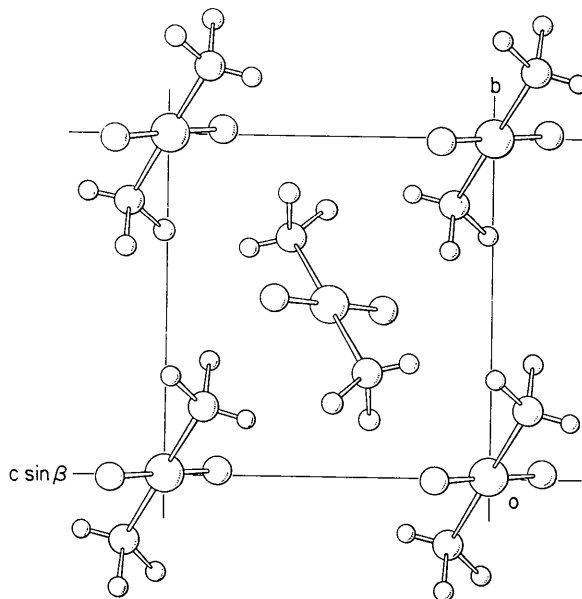
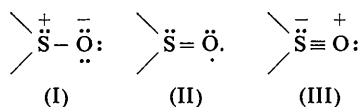


Fig. 2. Molecular packing in $\text{PdCl}_2(\text{DMSO})_2$ as shown by a projection down the crystallographic a axis.

* \bar{S}_2 computed assuming an isotropic data set containing 1300 independent reflections.

Table 6. Root mean square amplitudes and direction cosines of the atomic vibrational ellipsoids

	$(u^2_{\text{max}})^{1/2}$			$(u^2_{\text{med}})^{1/2}$			$(u^2_{\text{min}})^{1/2}$		
Pd		0.168			0.152			0.137	
S	0.918	0.084, -0.388		0.037, 0.955, 0.296			0.396, -0.286, 0.873		
		0.188			0.154			0.135	
Cl	0.982, -0.092, 0.096			0.059, 0.976, 0.213			-0.182, -0.199, 0.946		
	-0.652, 0.256	0.714		0.277, 0.957, -0.085			0.706, -0.143, 0.695		
O		0.276			0.204			0.145	
C(1)	0.945, -0.156, 0.289			0.205, 0.968, -0.144			-0.255, 0.196, 0.946		
		0.236			0.197			0.166	
C(2)	0.758, -0.639, -0.134			-0.052, 0.264, 0.962			0.651, 0.723, 0.235		
		0.244			0.199			0.149	
	0.934, 0.357, -0.031			-0.151, 0.314, 0.938			-0.324, 0.880, -0.347		



In (III) it must be assumed that sulfur is using at least one *d* orbital for back-accepting π electrons from oxygen.

Clearly, if (III) represents a reasonable electronic distribution but one which is far from predominant in DMSO itself, the attachment of sulfur to an electrophile should increase the contribution of (III) and consequently decrease the S–O bond length. In $(\text{CH}_3)_3\text{SO}^+$, where the very powerful electrophile CH_3^+ has been attached to the sulfur of DMSO, the S–O bond length has decreased to 1.45(1) Å, from the 1.513(5) Å found in DMSO. In $\text{PdCl}_2 \cdot (\text{DMSO})_2$, the S–O distance 1.475(5) Å is also shorter than in DMSO by 0.038(10) Å.

These changes in the S–O distance appear to be quite consistent with the corresponding changes in the S–O stretching frequency. This varies from 1096 cm^{-1} in gaseous DMSO to 1116 cm^{-1} in $\text{PdCl}_2(\text{DMSO})_2$ to 1233 cm^{-1} in $(\text{CH}_3)_3\text{SO}^+$.

The changes in bond length also seem to correlate in a reasonably semi-quantitative way with the assumed bond orders. Thus for S–N bonds one has a complete range of bond orders. Essentially a triple bond, with a length of 1.416(3) Å is found in F_3SN (Cotton & Wilkinson (1966) give a summary and references for the S–N bond lengths), while the S–N double bond in S_4N_4 has a length of 1.54(1) Å and the single bonds

found in $\text{S}_4\text{N}_4\text{H}_4$ have lengths of 1.66(1) and 1.67(1) Å respectively. Thus, for S–N bonds, the change in length per unit change in bond order is about 0.12 Å and a roughly similar value might be expected for S–O bonds. Moreover, the S–O bond of any given order ought to be a little shorter than the corresponding S–N bond due to the inherently smaller radii of O *versus* N.

Thus the S–O bond length of about 1.51 Å is consistent with the assignment of a bond order of about 2.0 when compared with the S–N bonds and the change from 1.51 to 1.45 Å on going from $(\text{CH}_3)_2\text{SO}$ to $(\text{CH}_3)_3\text{SO}^+$, that is 0.06 Å, would suggest that a very considerable increase, roughly to a bond order of $2 + (0.06)/(0.12) = 2.5$ had occurred. In $\text{PdCl}_2(\text{DMSO})_2$, on the other hand, a smaller change, 0.04 Å, perhaps to about 2.2, appears to have occurred.

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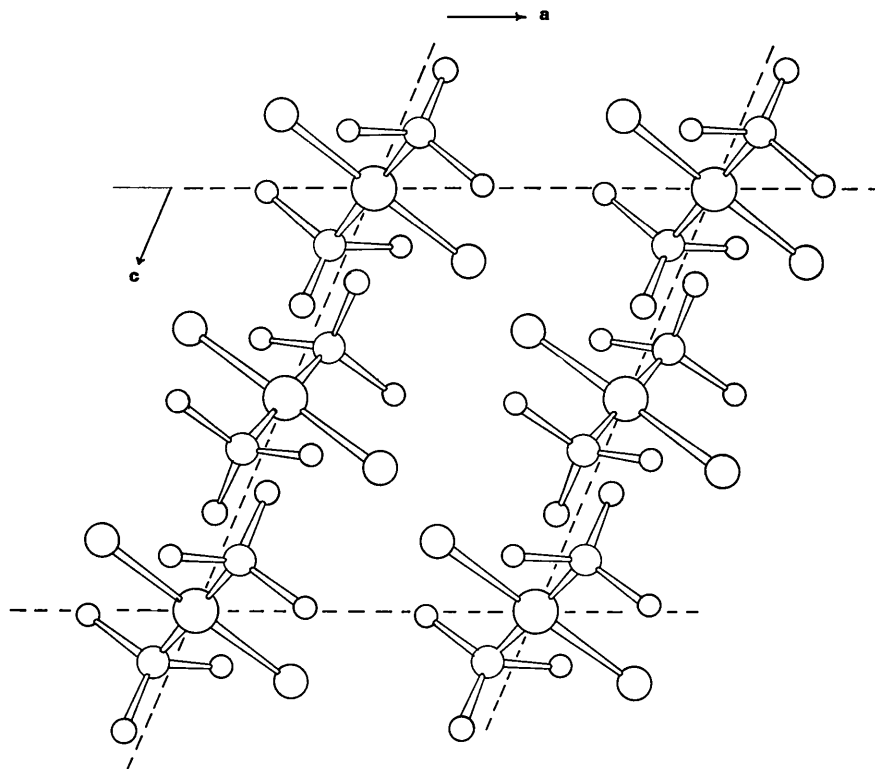


Fig. 3. Molecular packing in $\text{PdCl}_2(\text{DMSO})_2$ as shown by a projection down the crystallographic *b* axis.

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The Crystal Structure of Two Modifications of Tetraethylammonium Triiodide, $(C_2H_5)_4NI_3$

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The crystal structures of two modifications, I and II, of $(C_2H_5)_4NI_3$ have been determined by three-dimensional X-ray methods. Modification I crystallizes in space group *Cmca* with $a = 14.207$, $b = 15.220$, $c = 14.061$ Å, e.s.d. 0.05%, $Z = 8$; II crystallizes in space group *Pnma* with $a = 14.552$, $b = 13.893$, $c = 15.156$ Å, e.s.d. 0.05%, $Z = 8$. Anisotropic least-squares refinement was carried out with 909 independent reflexions for I ($R = 0.045$) and with 1094 reflexions for II ($R = 0.042$). In I, there are two independent I_3^- ions at special positions with symmetry $2/m$; the bond lengths in the two symmetrical I_3^- ions are 2.928 and 2.943 Å, e.s.d. 0.0025 Å. In II there are two independent I_3^- ions; the ions are nearly linear, but asymmetric with bond lengths of 2.912, 2.961 and 2.892, 2.981 Å, e.s.d. 0.0035 Å. A VESCF molecular orbital study of the electronic structure of the triiodide ions, including the effect of the crystal field, was made. A bond order *versus* bond length curve could be obtained.

Introduction

The nearly linear triiodide ion varies in structure with its environment in the crystal. Symmetric ions with equal bond lengths have, for instance, been observed in $(C_6H_5)_4AsI_3$ (Mooney Slater, 1959) whereas in crystals of CsI_3 (Tasman & Boswijk, 1955), the ions are asymmetric with a difference in I-I bond length of approximately 0.2 Å. From the data available at that time, Rundle (1961) concluded that the asymmetry and the total length of the ions increase with decreasing size of the surrounding cations. Hach & Rundle (1951) as well as Havinga & Wiebenga (1959) ascribe the variations in bond length to the influence of the (electrostatic) crystal field. Their qualitative conclusions

have recently been supported by Brown & Nunn (1966) for the triiodide ions in $(C_6H_5)_4AsI_3$ and in CsI_3 by theoretical variable electronegativity self-consistent field (VESCF) molecular orbital calculations in which the crystal field was taken into account. Mooney Slater (1959) and Slater (1959), on the other hand, have suggested that in certain crystals 'pressure' may force the I_3^- ions to become shorter and more symmetric.

Examination of the structural data on which these discussions are based shows that most of the data used do not give reliable information on the lengths of the I-I bonds in I_3^- ions. In $(CH_3)_4NI_5$, $(C_2H_5)_4NI_7$ and Cs_2I_8 , the I_3^- ions show strong interaction either with each other or with iodine molecules so that complexes with I_5^- , I_7^- and I_8^{2-} ions respectively occur in these